

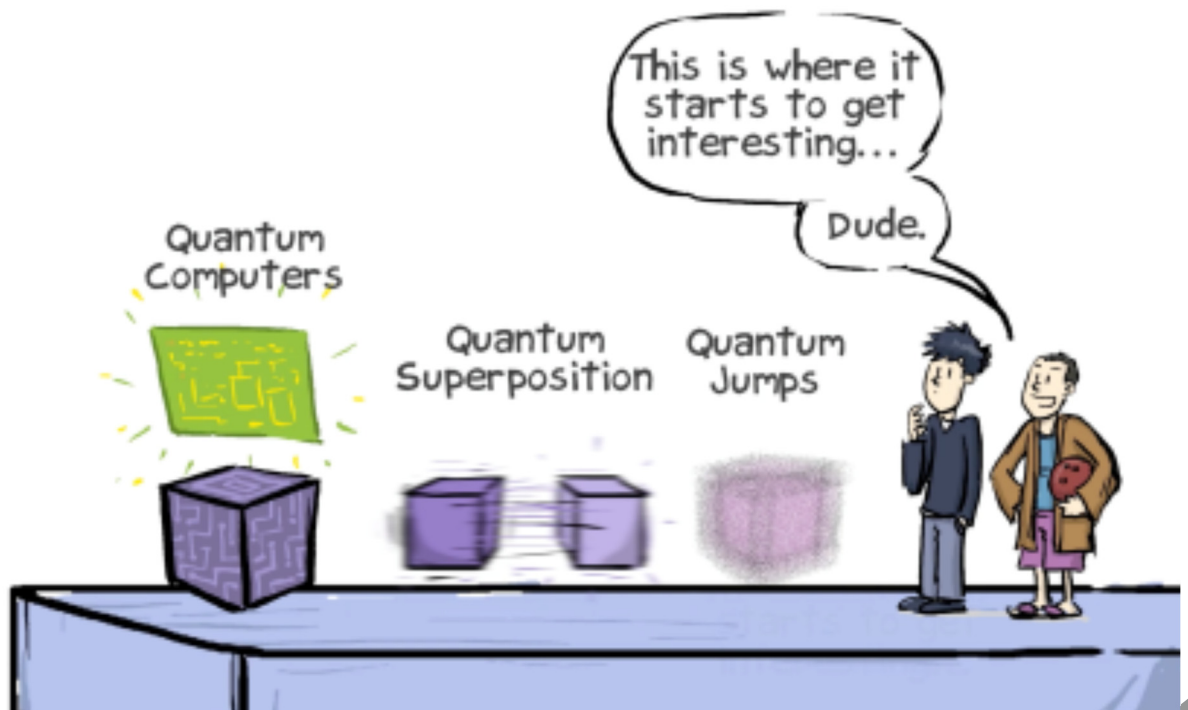
## Lecture 5

# Mean or Expectation Values of Observables in Quantum Mechanics

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In this lecture you will learn:

- How to obtain mean values of various physical quantities (observables) from wavefunctions



## The Schrödinger Equation: Recap

$$i\hbar \frac{\partial \psi(\mathbf{x}, t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi(\mathbf{x}, t)}{\partial \mathbf{x}^2} + V(\mathbf{x})\psi(\mathbf{x}, t)$$

When the potential is zero everywhere in space:

$$i\hbar \frac{\partial \psi(\mathbf{x}, t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi(\mathbf{x}, t)}{\partial \mathbf{x}^2}$$

The solution is a plane wave:

$$\psi(\mathbf{x}, t) = A e^{ikx} e^{-i\frac{E}{\hbar}t} \quad \text{Energy of the particle: } E(k) = \frac{\hbar^2 k^2}{2m}$$

Imagine the particle is in a very large box of length  $L$ :

$$\int_{-\infty}^{\infty} dx |\psi(\mathbf{x}, t)|^2 = \int_0^L dx |\psi(\mathbf{x}, t)|^2 = |A|^2 L = 1 \quad \Rightarrow \quad |A| = \frac{1}{\sqrt{L}}$$

Quantum wavefunction is defined only up to an overall phase factor:

$$A = \frac{e^{i\phi}}{\sqrt{L}}$$

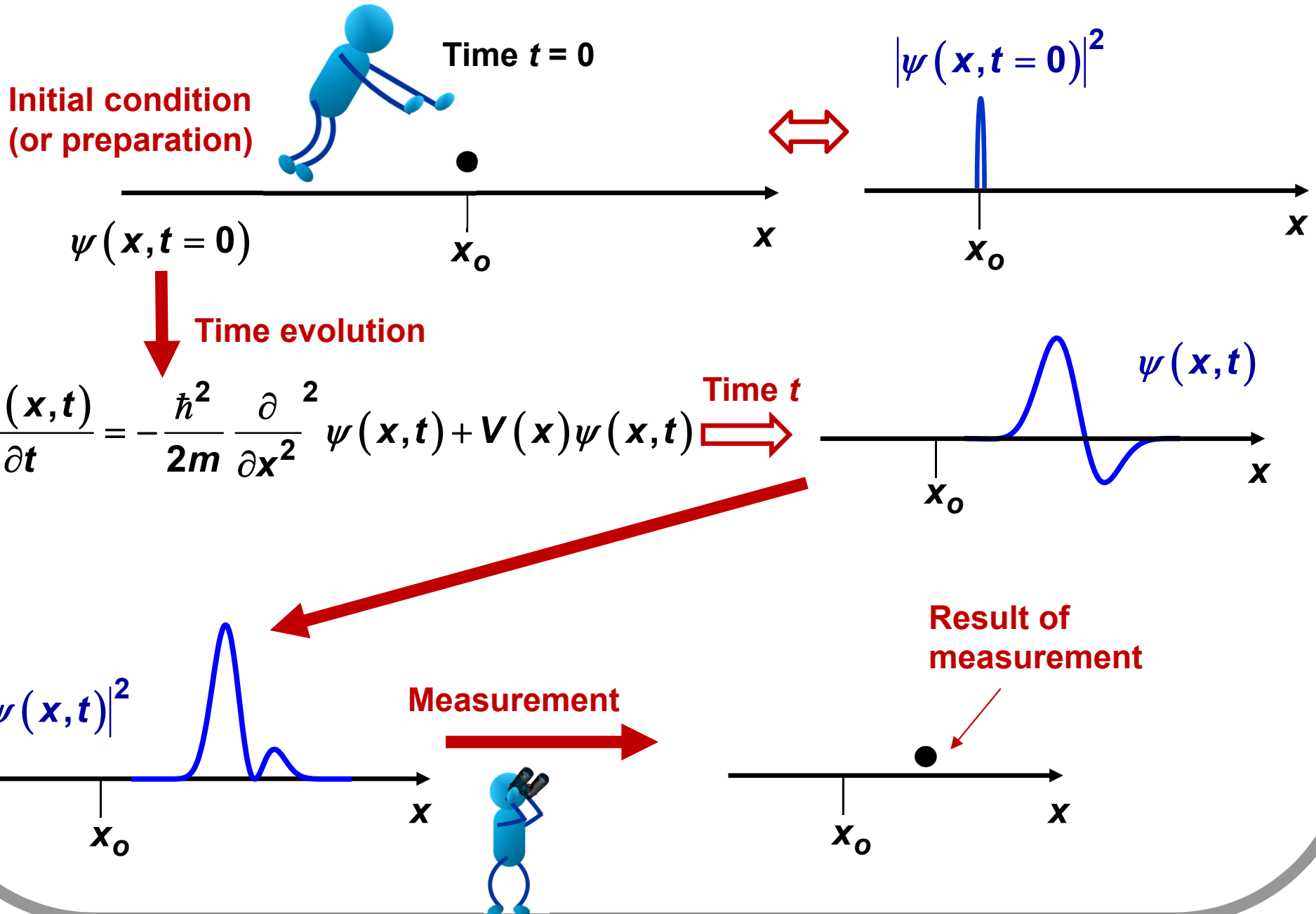
## Observables in Quantum Physics

An observable in quantum physics is any physical quantity that can be measured

Examples of observables associated with a particle include its position, its momentum, its energy, its angular momentum, its spin, etc

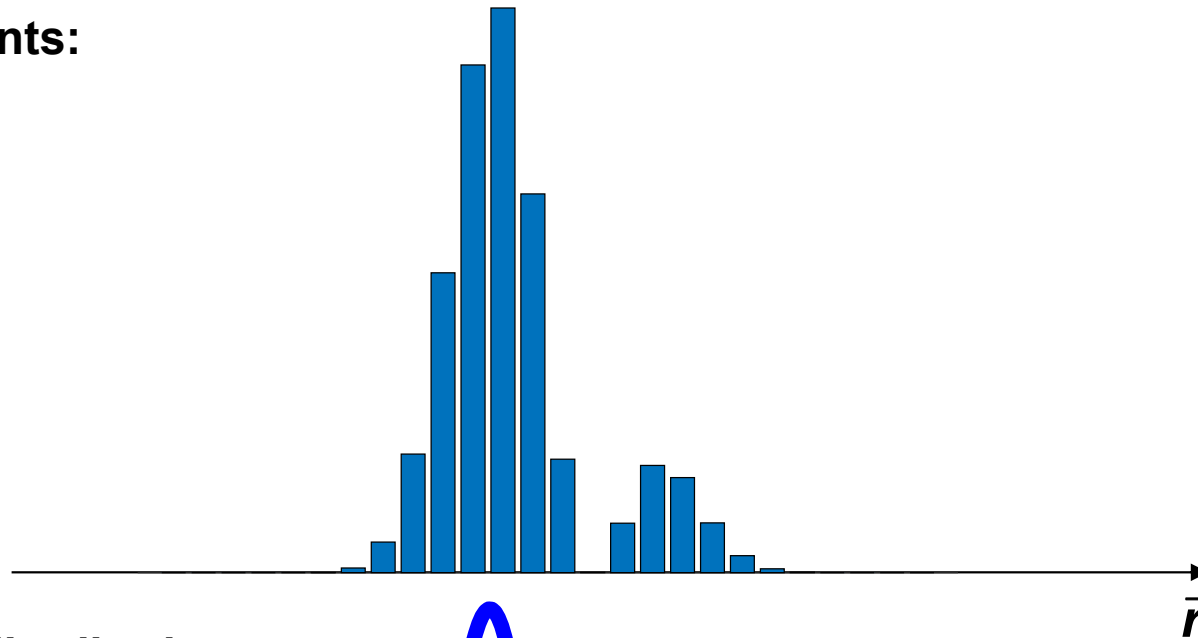
**Question:** How do we obtain the mean values of various physical quantities (observables) from the wavefunction?

# The Statistical Interpretation of Quantum Probabilities

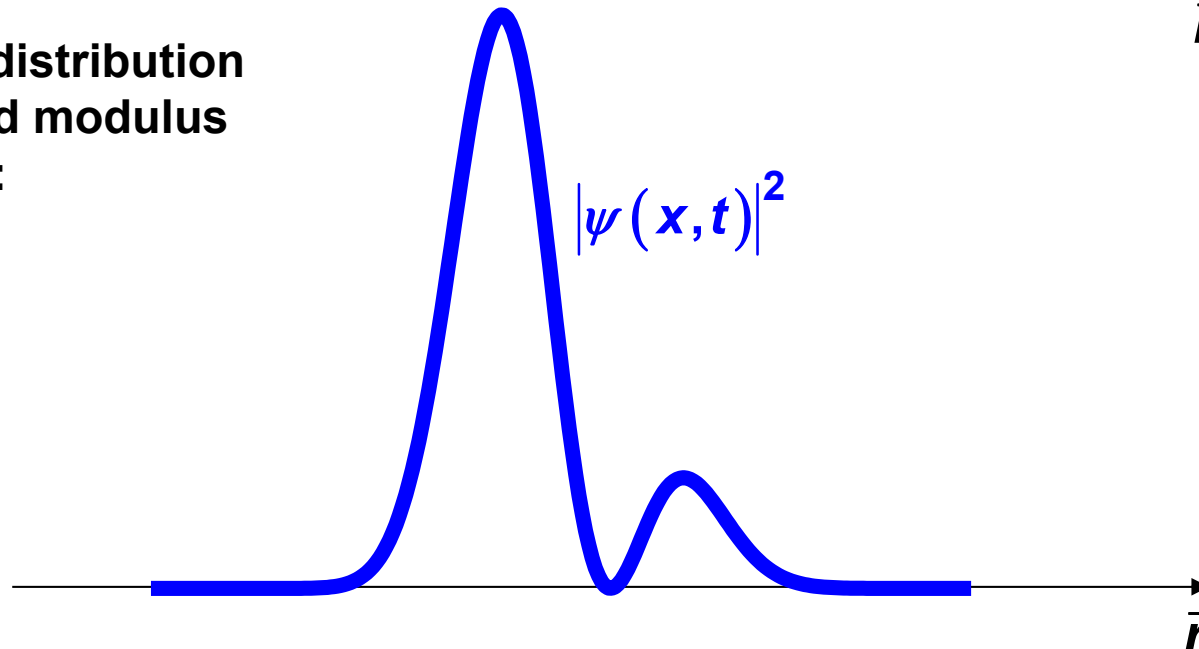


# The Statistical Interpretation of Quantum Probabilities

Histogram of the results of position measurements:



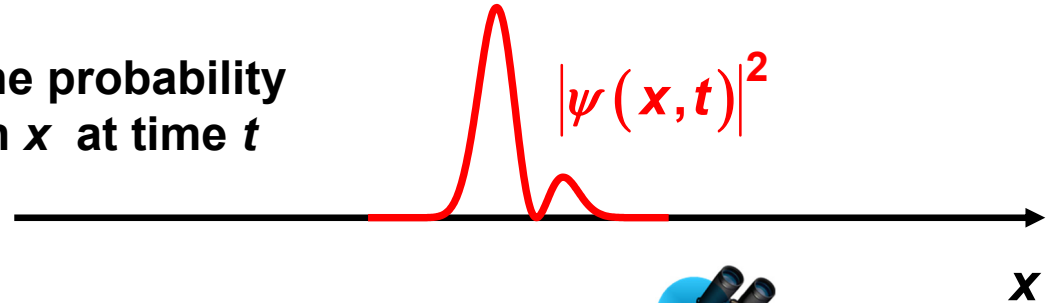
A-priori probability distribution given by the squared modulus of the wavefunction:



## Mean Value of Position or Position Expectation Value

Since the a-priori probability (or rather the probability density) of finding the particle at location  $x$  at time  $t$  is:

$$|\psi(x, t)|^2$$



The **average position** of the particle at time  $t$  is then:

$$\langle x \rangle(t) = \int_{-\infty}^{\infty} dx \ x |\psi(x, t)|^2 = \int_{-\infty}^{\infty} dx \ \psi^*(x, t) \underbrace{[x]}_{\text{Position operator}} \psi(x, t)$$



This is also called the **expectation value** of the position, or the **mean value** of the position, or the **average value** of the position

**Meaning:** If the experiment is repeated many times under identical conditions, and each time a measurement is made at time  $t$  to locate the particle, then the mean location is determined experimentally by taking the average of the measurement results obtained in all these experiments and this experimental mean would equal the **expectation value** or the **mean value** given above by the quantum wavefunction

# Mean Value of Momentum or Momentum Expectation Value

We know that in classical physics:

$$p(t) = mv(t) = m \frac{dx(t)}{dt}$$

So we expect the **momentum expectation value** in quantum mechanics to go as:

$$\langle p \rangle(t) = m \frac{d}{dt} \langle x \rangle(t)$$

$$i\hbar \frac{\partial \psi(x,t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x,t)}{\partial x^2} + V(x)\psi(x,t)$$

$$-i\hbar \frac{\partial \psi^*(x,t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi^*(x,t)}{\partial x^2} + V(x)\psi^*(x,t)$$

We start from:

$$\langle x \rangle(t) = \int_{-\infty}^{\infty} dx \psi^*(x,t) [x] \psi(x,t)$$

$$\Rightarrow m \frac{d}{dt} \langle x \rangle(t) = \int_{-\infty}^{\infty} dx m \frac{\partial \psi^*(x,t)}{\partial t} [x] \psi(x,t) + \int_{-\infty}^{\infty} dx \psi^*(x,t) [x] m \frac{\partial \psi(x,t)}{\partial t}$$

use

use

# Mean Value of Momentum or Momentum Expectation Value

$$\langle x \rangle(t) = \int_{-\infty}^{\infty} dx \psi^*(x,t)[x]\psi(x,t)$$

$$\Rightarrow m \frac{d}{dt} \langle x \rangle(t) = \int_{-\infty}^{\infty} dx m \frac{\partial \psi^*(x,t)}{\partial t} [x] \psi(x,t) + \int_{-\infty}^{\infty} dx \psi^*(x,t) [x] m \frac{\partial \psi(x,t)}{\partial t}$$

$$= \int_{-\infty}^{\infty} dx \left( \frac{\hbar}{2i} \right) \frac{\partial^2 \psi^*(x,t)}{\partial x^2} [x] \psi(x,t) + \int_{-\infty}^{\infty} dx \psi^*(x,t) [x] \left( -\frac{\hbar}{2i} \right) \frac{\partial^2 \psi(x,t)}{\partial x^2}$$

Integrate by parts both terms

Integrate by parts first term

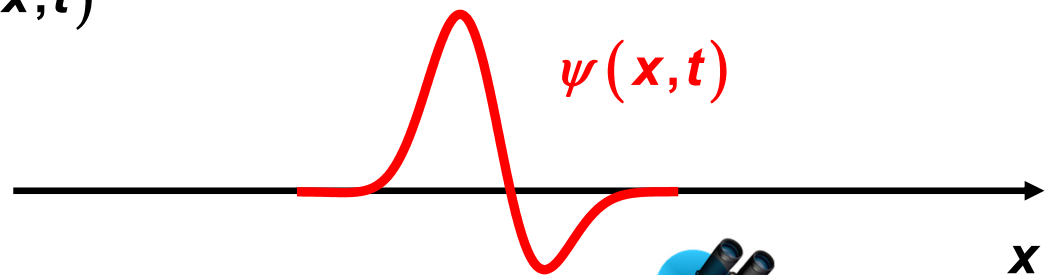
$$= \int_{-\infty}^{\infty} dx \left( -\frac{\hbar}{2i} \right) \frac{\partial \psi^*(x,t)}{\partial x} \psi(x,t) + \int_{-\infty}^{\infty} dx \psi^*(x,t) \left( \frac{\hbar}{2i} \right) \frac{\partial \psi(x,t)}{\partial x}$$

$$= \int_{-\infty}^{\infty} dx \psi^*(x,t) \left[ \frac{\hbar}{i} \frac{\partial}{\partial x} \right] \psi(x,t)$$

The final answer becomes:

$$\langle p \rangle(t) = \int_{-\infty}^{\infty} dx \psi^*(x,t) \left[ \frac{\hbar}{i} \frac{\partial}{\partial x} \right] \psi(x,t)$$

Momentum operator



Measurement



## Mean Value of Kinetic Energy (KE) or KE Expectation Value

We know that in classical physics:

$$\text{KE} = \frac{p^2(t)}{2m}$$

So we expect the **KE expectation value** in quantum mechanics to go as:

$$\langle \text{KE} \rangle(t) = \frac{\langle p^2 \rangle(t)}{2m}$$

Since:

$$\langle p \rangle(t) = \int_{-\infty}^{\infty} dx \psi^*(x,t) \left[ \frac{\hbar}{i} \frac{\partial}{\partial x} \right] \psi(x,t)$$

We conclude that:

$$\begin{aligned} \langle \text{KE} \rangle(t) &= \frac{\langle p^2 \rangle(t)}{2m} = \frac{1}{2m} \int_{-\infty}^{\infty} dx \psi^*(x,t) \left[ \frac{\hbar}{i} \frac{\partial}{\partial x} \frac{\hbar}{i} \frac{\partial}{\partial x} \right] \psi(x,t) \\ &= \int_{-\infty}^{\infty} dx \psi^*(x,t) \underbrace{\left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \right]}_{\text{KE operator}} \psi(x,t) \end{aligned}$$

## Mean Value of Potential Energy (PE) or PE Expectation Value

Since the a-priori probability of finding the particle at location  $x$  at time  $t$  is:

$$|\psi(x, t)|^2$$

And the **average position** of the particle at time  $t$  is:

$$\langle x \rangle(t) = \int_{-\infty}^{\infty} dx \psi^*(x, t) [x] \psi(x, t)$$

We expect:

$$\langle V(x) \rangle(t) = \int_{-\infty}^{\infty} dx \psi^*(x, t) \underbrace{[V(x)]}_{\text{PE operator}} \psi(x, t)$$

PE operator

## Mean Value of Energy or Energy Expectation Value

The total energy of a particle is the sum of its kinetic and potential energies

So we expect that in quantum mechanics the mean value of the total energy will be:

$$\begin{aligned}\langle E \rangle(t) &= \langle KE \rangle(t) + \langle PE \rangle(t) = \frac{\langle p^2 \rangle(t)}{2m} + \langle V(x) \rangle(t) \\ &= \int_{-\infty}^{\infty} dx \psi^*(x,t) \left[ \underbrace{-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x)}_{\text{Total energy operator}} \right] \psi(x,t)\end{aligned}$$

Another way of writing the Schrodinger equation:

$$\begin{aligned}i\hbar \frac{\partial \psi(x,t)}{\partial t} &= -\frac{\hbar^2}{2m} \frac{\partial^2 \psi(x,t)}{\partial x^2} + V(x) \psi(x,t) \\ \Rightarrow i\hbar \frac{\partial \psi(x,t)}{\partial t} &= \left[ \text{Energy operator} \right] \psi(x,t)\end{aligned}$$

## Mean Values or Expectation Values of Observables in 1D

Position:

$$\langle \mathbf{x} \rangle(t) = \int_{-\infty}^{\infty} d\mathbf{x} \psi^*(\mathbf{x}, t) [\mathbf{x}] \psi(\mathbf{x}, t)$$

Potential Energy:

$$\langle \mathbf{PE} \rangle(t) = \langle \mathbf{V}(\mathbf{x}) \rangle(t) = \int_{-\infty}^{\infty} d\mathbf{x} \psi^*(\mathbf{x}, t) [\mathbf{V}(\mathbf{x})] \psi(\mathbf{x}, t)$$

Momentum:

$$\langle \mathbf{p} \rangle(t) = \int_{-\infty}^{\infty} d\mathbf{x} \psi^*(\mathbf{x}, t) \left[ \frac{\hbar}{i} \frac{\partial}{\partial \mathbf{x}} \right] \psi(\mathbf{x}, t)$$

Kinetic Energy:

$$\langle \mathbf{KE} \rangle(t) = \int_{-\infty}^{\infty} d\mathbf{x} \psi^*(\mathbf{x}, t) \left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial \mathbf{x}^2} \right] \psi(\mathbf{x}, t)$$

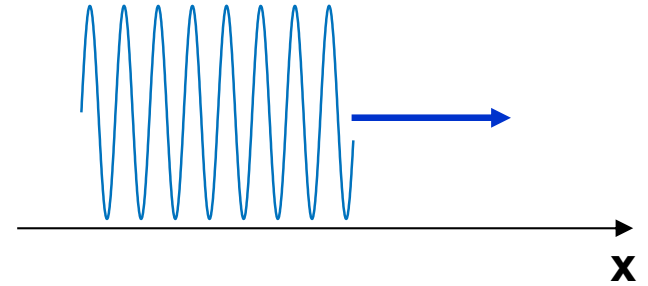
Total Energy:

$$\langle \mathbf{E} \rangle(t) = \int_{-\infty}^{\infty} d\mathbf{x} \psi^*(\mathbf{x}, t) \left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial \mathbf{x}^2} + \mathbf{V}(\mathbf{x}) \right] \psi(\mathbf{x}, t)$$

## Mean Values or Expectation Value in 1D: Example

Consider a particle in 1D in **free space** (potential is zero everywhere):

$$\psi(x, t) = \frac{1}{\sqrt{L}} e^{ikx} e^{-i\frac{E(k)}{\hbar}t}$$



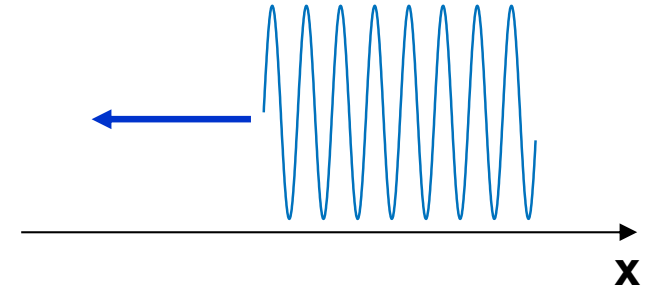
Mean value of momentum:

$$\begin{aligned}\langle p \rangle(t) &= \int_{-\infty}^{\infty} dx \psi^*(x, t) \left[ \frac{\hbar}{i} \frac{\partial}{\partial x} \right] \psi(x, t) \\ &= \int_{-\infty}^{\infty} dx \frac{1}{\sqrt{L}} e^{-ikx} e^{+i\frac{E(k)}{\hbar}t} \left[ \frac{\hbar}{i} \frac{\partial}{\partial x} \right] \frac{1}{\sqrt{L}} e^{ikx} e^{-i\frac{E(k)}{\hbar}t} \\ &= \int_{-\infty}^{\infty} dx \frac{1}{\sqrt{L}} e^{-ikx} e^{+i\frac{E(k)}{\hbar}t} [\hbar k] \frac{1}{\sqrt{L}} e^{ikx} e^{-i\frac{E(k)}{\hbar}t} \\ &= \hbar k \int_{-\infty}^{\infty} dx \frac{1}{\sqrt{L}} e^{-ikx} e^{+i\frac{E(k)}{\hbar}t} \frac{1}{\sqrt{L}} e^{ikx} e^{-i\frac{E(k)}{\hbar}t} \\ &= \hbar k\end{aligned}$$

## Mean Values or Expectation Value in 1D: Example

Consider a particle in 1D in **free space** (potential is zero everywhere):

$$\psi(\mathbf{x}, t) = \frac{1}{\sqrt{L}} e^{-ikx} e^{-i\frac{E(-k)}{\hbar}t} = \frac{1}{\sqrt{L}} e^{-ikx} e^{-i\frac{E(k)}{\hbar}t}$$



Mean value of momentum:

$$\begin{aligned} \langle p \rangle(t) &= \int_{-\infty}^{\infty} d\mathbf{x} \psi^*(\mathbf{x}, t) \left[ \frac{\hbar}{i} \frac{\partial}{\partial \mathbf{x}} \right] \psi(\mathbf{x}, t) \\ &= \int_{-\infty}^{\infty} d\mathbf{x} \frac{1}{\sqrt{L}} e^{+ikx} e^{+i\frac{E(k)}{\hbar}t} \left[ \frac{\hbar}{i} \frac{\partial}{\partial \mathbf{x}} \right] \frac{1}{\sqrt{L}} e^{-ikx} e^{-i\frac{E(k)}{\hbar}t} \\ &= \int_{-\infty}^{\infty} d\mathbf{x} \frac{1}{\sqrt{L}} e^{+ikx} e^{+i\frac{E(k)}{\hbar}t} [-\hbar k] \frac{1}{\sqrt{L}} e^{-ikx} e^{-i\frac{E(k)}{\hbar}t} \\ &= -\hbar k \int_{-\infty}^{\infty} d\mathbf{x} \frac{1}{\sqrt{L}} e^{+ikx} e^{+i\frac{E(k)}{\hbar}t} \frac{1}{\sqrt{L}} e^{-ikx} e^{-i\frac{E(k)}{\hbar}t} \\ &= -\hbar k \end{aligned}$$

## Mean Values or Expectation Value in 1D: Example

Consider a particle in 1D in **free space** (potential is zero everywhere):

$$\begin{aligned}\psi(x, t) &= \frac{1}{\sqrt{2L}} e^{ikx} e^{-i\frac{E(k)}{\hbar}t} + \frac{1}{\sqrt{2L}} e^{-ikx} e^{-i\frac{E(-k)}{\hbar}t} \\ &= \sqrt{\frac{2}{L}} \cos(kL) e^{-i\frac{E(k)}{\hbar}t}\end{aligned}$$

The state is an equal-weight superposition of two plane waves

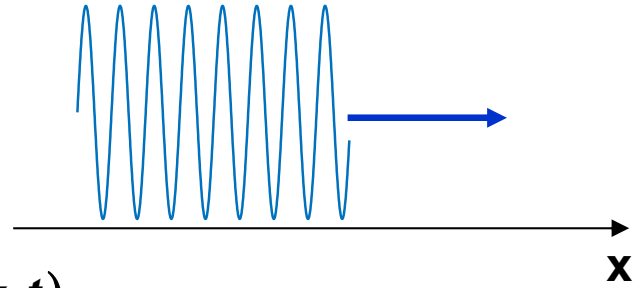
**Mean value of momentum:**

$$\begin{aligned}\langle p \rangle(t) &= \int_{-\infty}^{\infty} dx \psi^*(x, t) \left[ \frac{\hbar}{i} \frac{\partial}{\partial x} \right] \psi(x, t) \\ &= 0\end{aligned}$$

## Mean Values or Expectation Value in 1D: Example

Consider a particle in 1D in **free space** (potential is zero everywhere):

$$\psi(x, t) = \frac{1}{\sqrt{L}} e^{ikx} e^{-i\frac{E(k)}{\hbar}t}$$



Mean  
energy:

$$\begin{aligned} \langle E \rangle(t) &= \int_{-\infty}^{\infty} dx \psi^*(x, t) \left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \cancel{V(x)} \right] \psi(x, t) \\ &= \int_{-\infty}^{\infty} dx \psi^*(x, t) \left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \right] \psi(x, t) \\ &= \int_{-\infty}^{\infty} dx \frac{1}{\sqrt{L}} e^{-ikx} e^{+i\frac{E(k)}{\hbar}t} \left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \right] \frac{1}{\sqrt{L}} e^{ikx} e^{-i\frac{E(k)}{\hbar}t} \\ &= \int_{-\infty}^{\infty} dx \frac{1}{\sqrt{L}} e^{-ikx} e^{+i\frac{E(k)}{\hbar}t} \left[ \frac{\hbar^2 k^2}{2m} \right] \frac{1}{\sqrt{L}} e^{ikx} e^{-i\frac{E(k)}{\hbar}t} \\ &= \frac{\hbar^2 k^2}{2m} \int_{-\infty}^{\infty} dx \frac{1}{\sqrt{L}} e^{-ikx} e^{+i\frac{E(k)}{\hbar}t} \frac{1}{\sqrt{L}} e^{ikx} e^{-i\frac{E(k)}{\hbar}t} \\ &= \frac{\hbar^2 k^2}{2m} \end{aligned}$$



## Fourier Transformed Wavefunctions

Consider a particle in 1D with wavefunction:  $\psi(x, t)$

The Fourier transform of the wavefunction is:

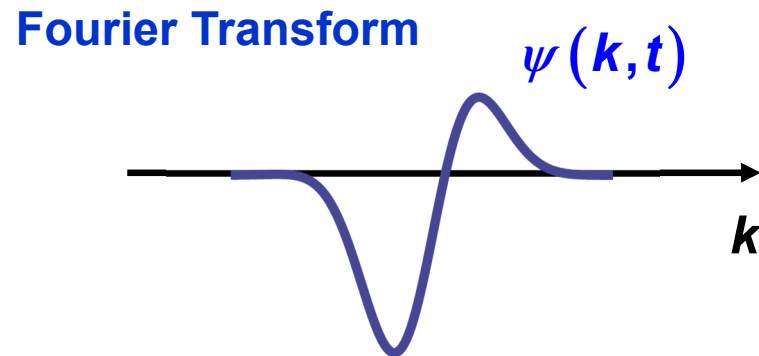
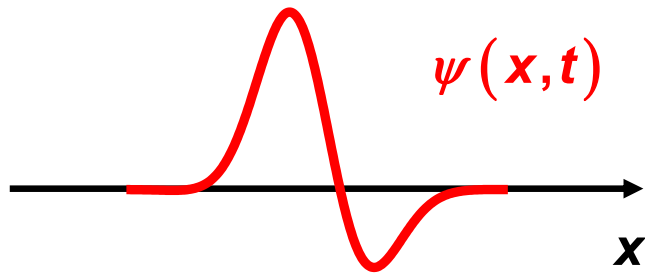
$$\psi(k, t) = \int_{-\infty}^{\infty} dx \psi(x, t) e^{-ikx}$$

$$\psi(x, t) = \int_{-\infty}^{\infty} \frac{dk}{2\pi} \psi(k, t) e^{ikx}$$

**Note: unless the particle is free do not assume:**

$$\psi(k, t) = \psi(k) e^{-\frac{i}{\hbar} E(k)t}$$

$$E(k) = \frac{\hbar^2 k^2}{2m}$$



And:

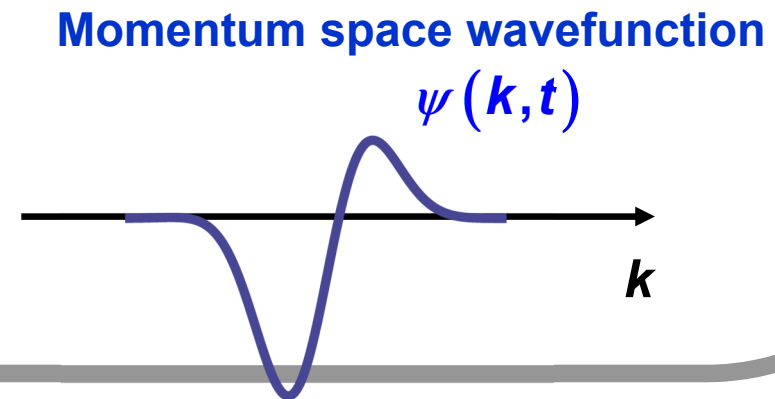
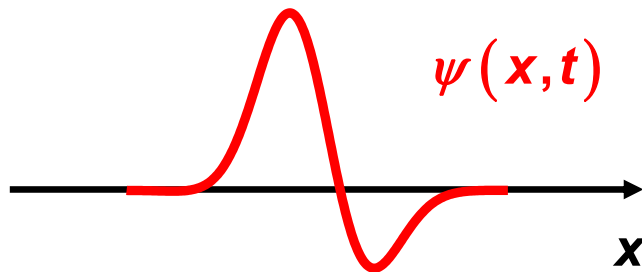
$$\int_{-\infty}^{\infty} dx |\psi(x, t)|^2 = 1 \Rightarrow \int_{-\infty}^{\infty} \frac{dk}{2\pi} |\psi(k, t)|^2 = 1$$

## Fourier Transformed Wavefunctions

$$\psi(k, t) = \int_{-\infty}^{\infty} dx \psi(x, t) e^{-ikx}$$

Momentum expectation value can also be expressed as:

$$\begin{aligned} \langle p \rangle(t) &= \int_{-\infty}^{\infty} dx \psi^*(x, t) \left[ \frac{\hbar}{i} \frac{\partial}{\partial x} \right] \psi(x, t) \\ &= \int_{-\infty}^{\infty} \frac{dk}{2\pi} \int_{-\infty}^{\infty} \frac{dk'}{2\pi} \int_{-\infty}^{\infty} dx \psi^*(k, t) e^{-ikx} [\hbar k'] \psi(k', t) e^{ik'x} \\ &= \int_{-\infty}^{\infty} \frac{dk}{2\pi} \int_{-\infty}^{\infty} \frac{dk'}{2\pi} \psi^*(k, t) [\hbar k'] \psi(k', t) \int_{-\infty}^{\infty} dx e^{-ikx} e^{ik'x} \\ &= \int_{-\infty}^{\infty} \frac{dk}{2\pi} \int_{-\infty}^{\infty} \frac{dk'}{2\pi} \hbar k' \psi^*(k, t) \psi(k', t) (2\pi) \delta(k - k') \\ &= \int_{-\infty}^{\infty} \frac{dk}{2\pi} \hbar k |\psi(k, t)|^2 \end{aligned}$$



## Gaussian Wavepacket Example

Consider a particle in 1D in **free space** (potential is zero everywhere):

$$\psi(\mathbf{x}, t) = \int_{-\infty}^{\infty} \frac{d\mathbf{k}}{2\pi} \psi(\mathbf{k}, t) e^{i\mathbf{k}\mathbf{x}}$$

“Wavefunction” in  
Fourier space

Normalization:

$$\int_{-\infty}^{\infty} d\mathbf{x} |\psi(\mathbf{x}, t)|^2 = 1 \quad \Rightarrow \quad \int_{-\infty}^{\infty} \frac{d\mathbf{k}}{2\pi} |\psi(\mathbf{k}, t)|^2 = 1$$

Free-space Shrodinger equation (zero potential):

$$i\hbar \frac{\partial \psi(\mathbf{x}, t)}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial \mathbf{x}^2} \psi(\mathbf{x}, t)$$

Substitute the Fourier expression:

$$\psi(\mathbf{x}, t) = \int_{-\infty}^{\infty} \frac{d\mathbf{k}}{2\pi} \psi(\mathbf{k}, t) e^{i\mathbf{k}\mathbf{x}}$$

To get:

$$\int_{-\infty}^{\infty} \frac{d\mathbf{k}}{2\pi} i\hbar \frac{\partial}{\partial t} \psi(\mathbf{k}, t) e^{i\mathbf{k}\mathbf{x}} = \int_{-\infty}^{\infty} \frac{d\mathbf{k}}{2\pi} \frac{\hbar^2 \mathbf{k}^2}{2m} \psi(\mathbf{k}, t) e^{i\mathbf{k}\mathbf{x}}$$

## Gaussian Wavepacket Example

$$\int_{-\infty}^{\infty} \frac{dk}{2\pi} i\hbar \frac{\partial}{\partial t} \psi(k, t) e^{ikx} = \int_{-\infty}^{\infty} \frac{dk}{2\pi} \frac{\hbar^2 k^2}{2m} \psi(k, t) e^{ikx}$$

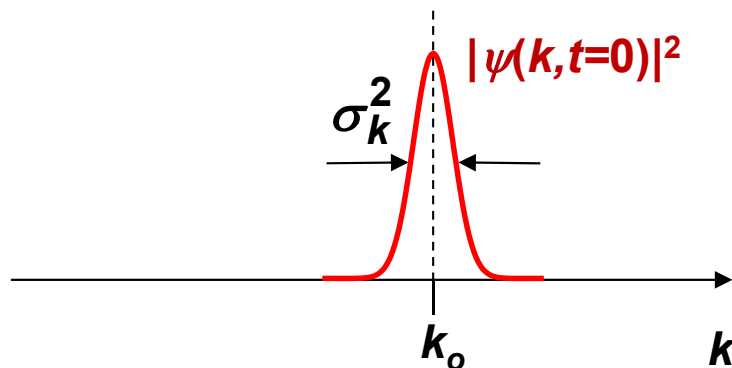
$$\Rightarrow i\hbar \frac{\partial}{\partial t} \psi(k, t) = E(k) \psi(k, t) \quad E(k)$$

$$\Rightarrow \psi(k, t) = \phi(k) e^{-\frac{i}{\hbar} E(k)t}$$

This implies:

$$\psi(x, t) = \int_{-\infty}^{\infty} \frac{dk}{2\pi} \psi(k, t) e^{ikx} = \int_{-\infty}^{\infty} \frac{dk}{2\pi} \phi(k) e^{ikx} e^{-\frac{i}{\hbar} E(k)t}$$

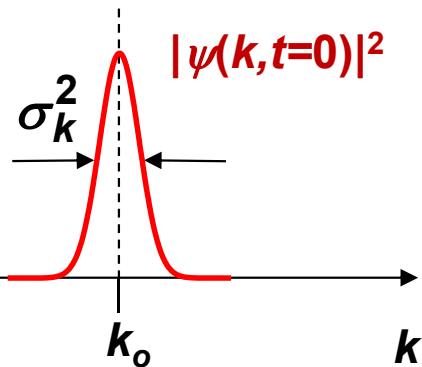
Consider a Gaussian packet in wavevector space or Fourier space:



$$\psi(k, t=0) = \phi(k) = \left( \frac{2\pi}{\sigma_k^2} \right)^{1/4} e^{-\frac{(k-k_0)^2}{4\sigma_k^2}}$$

## Gaussian Wavepacket Example

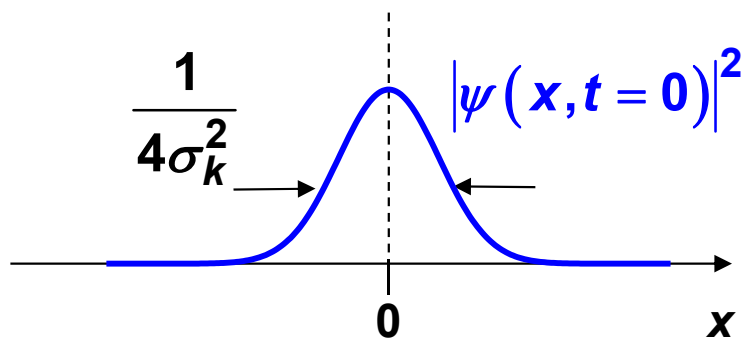
In wavevector or Fourier space:



$$\psi(k, t) = \phi(k) e^{-i \frac{E(k)}{\hbar} t} = \left( \frac{2\pi}{\sigma_k^2} \right)^{1/4} e^{-\frac{(k-k_0)^2}{4\sigma_k^2}} e^{-i \frac{E(k)}{\hbar} t}$$

$$\psi(k, t=0) = \phi(k) = \left( \frac{2\pi}{\sigma_k^2} \right)^{1/4} e^{-\frac{(k-k_0)^2}{4\sigma_k^2}}$$

In real space:



$$\psi(x, t) = \int_{-\infty}^{\infty} \frac{dk}{2\pi} \left[ \left( \frac{2\pi}{\sigma_k^2} \right)^{1/4} e^{-\frac{(k-k_0)^2}{4\sigma_k^2}} \right] e^{ikx} e^{-i \frac{E(k)}{\hbar} t}$$

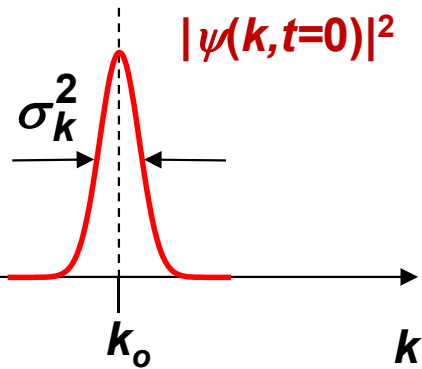
$$\Rightarrow \psi(x, t=0) = \left( \frac{4\sigma_k^2}{2\pi} \right)^{1/4} e^{-\frac{x^2}{4(1/4\sigma_k^2)}} e^{ik_0 x}$$

Nobel prize winning statement

A fat wavepacket in real space requires a small number of superposed plane waves to make it  
 A slim wavepacket in real space requires a large number of superposed plane waves to make it

## Gaussian Wavepacket Example

In Fourier space:



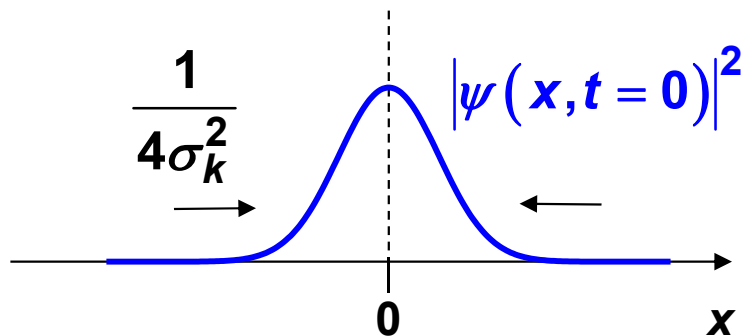
$$\psi(k, t) = \left( \frac{2\pi}{\sigma_k^2} \right)^{1/4} e^{-\frac{(k-k_0)^2}{4\sigma_k^2}} e^{-i\frac{E(k)}{\hbar}t}$$

$$\psi(k, t=0) = \left( \frac{2\pi}{\sigma_k^2} \right)^{1/4} e^{-\frac{(k-k_0)^2}{4\sigma_k^2}}$$

Width of the packet in Fourier space:

$$\sigma_k^2 = \int_{-\infty}^{\infty} \frac{dk}{2\pi} (k - k_0)^2 |\psi(k, t=0)|^2$$

In real space:



$$\psi(x, t) = \int_{-\infty}^{\infty} \frac{dk}{2\pi} \psi(k, t) e^{ikx}$$

$$\psi(x, t=0) = \left( \frac{4\sigma_k^2}{2\pi} \right)^{1/4} e^{-\frac{x^2}{4(1/4\sigma_k^2)}} e^{ik_0x}$$

Width of the packet in real space:

$$\sigma_x^2 = \int_{-\infty}^{\infty} dx x^2 |\psi(x, t=0)|^2 = \frac{1}{4\sigma_k^2}$$

$$\sigma_x^2 \sigma_k^2 = \frac{1}{4}$$

## Gaussian Wavepacket Example

Now momentum  $p$  equals  $\hbar k$

Therefore, the standard deviation of momentum at time  $t=0$  is:

$$\sigma_p^2 = \int_{-\infty}^{\infty} \frac{dk}{2\pi} (\hbar k - \hbar k_0)^2 |\psi(k, t=0)|^2 = \hbar^2 \sigma_k^2$$

The standard deviation of position was:

$$\sigma_x^2 = \int_{-\infty}^{\infty} dx x^2 |\psi(x, t=0)|^2 = \frac{1}{4\sigma_k^2}$$

This implies:

$$\sigma_x^2 \sigma_p^2 = \frac{\hbar^2}{4}$$

or

$$\sigma_x \sigma_p = \frac{\hbar}{2}$$

A fat wavepacket in real space requires a small number of superposed plane waves to make it

A slim wavepacket in real space requires a large number of superposed plane waves to make it

## Mean Values or Expectation Values in 3D

**Position:**

$$\langle \vec{r} \rangle(t) = \int d^3\vec{r} \psi^*(\vec{r}, t) [\vec{r}] \psi(\vec{r}, t)$$

**Potential Energy:**

$$\langle \text{PE} \rangle(t) = \langle V(\vec{r}) \rangle(t) = \int d^3\vec{r} \psi^*(\vec{r}, t) [V(\vec{r})] \psi(\vec{r}, t)$$

**Momentum:**

$$\langle \vec{p} \rangle(t) = \int d^3\vec{r} \psi^*(\vec{r}, t) \left[ \frac{\hbar}{i} \nabla \right] \psi(\vec{r}, t)$$

**Kinetic Energy:**

$$\langle \text{KE} \rangle(t) = \left\langle \frac{p^2}{2m} \right\rangle(t) = \int d^3\vec{r} \psi^*(\vec{r}, t) \left[ -\frac{\hbar^2}{2m} \nabla^2 \right] \psi(\vec{r}, t)$$

**Total Energy:**

$$\langle E \rangle(t) = \int d^3\vec{r} \psi^*(\vec{r}, t) \left[ -\frac{\hbar^2}{2m} \nabla^2 + V(\vec{r}) \right] \psi(\vec{r}, t)$$



