

Lecture 8

Boundary Value Problems

In this lecture you will learn:

- How to solve some interesting boundary value problems

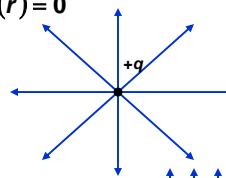
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Solutions of Laplace Equation in Spherical Coordinates

For all the following solutions $\nabla^2 \phi(\vec{r}) = 0$

Spherically Symmetric Solution

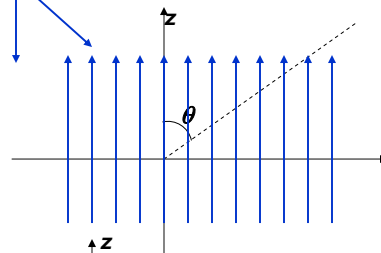
$$\phi(\vec{r}) = \frac{A}{r} + B$$



A Constant Uniform Electric Field Solution

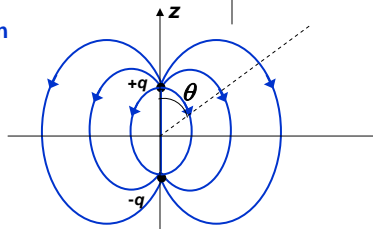
$$\phi(\vec{r}) = -A r \cos(\theta) = -Az$$

$$\vec{E}(\vec{r}) = -\nabla \phi(\vec{r}) = A \hat{z}$$



A Dipole Oriented Along z-Axis Solution

$$\phi(\vec{r}) = A \frac{\cos(\theta)}{r^2}$$

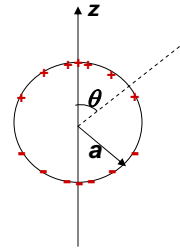


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
A Non-Uniformly Charged Spherical Shell - I

Consider a charged spherical sheet where the surface charge density is fixed and is given by:


$$\sigma = \sigma_0 \cos(\theta)$$




Looking from outside the potential sort of looks like that of a dipole. So try a **dipole-like solution** for $r > a$

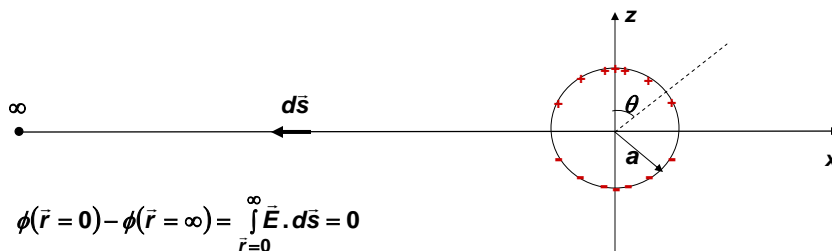
For $r > a$: $\phi_{out}(\vec{r}) = A \frac{\cos(\theta)}{r^2} + C$ 
Why??

For $r < a$: Try something that does not go to infinity at $r = 0$

$$\phi_{in}(\vec{r}) = B r \cos(\theta) + D$$
 
Why??

A Non-Uniformly Charged Spherical Shell - II

For $r < a$: $\phi_{in}(\vec{r}) = B r \cos(\theta) + D$ 
Why??



$$\phi(\vec{r} = 0) - \phi(\vec{r} = \infty) = \int_{\vec{r}=0}^{\infty} \vec{E} \cdot d\vec{s} = 0$$

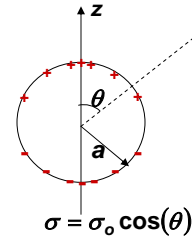
How would you know if your solution is the right one??

Uniqueness Theorem:

There is only one right solution and if you have a solution that satisfies all the boundary conditions then you have the right solution

A Non-Uniformly Charged Spherical Shell - III

$$\phi_{out}(\vec{r}) = A \frac{\cos(\theta)}{r^2} \quad \phi_{in}(\vec{r}) = B r \cos(\theta)$$



Boundary Conditions:

- (1) Potential inside and outside must be continuous at the surface of the shell

$$\begin{aligned} \phi_{in}(\vec{r})|_{r=a} &= \phi_{out}(\vec{r})|_{r=a} \\ \Rightarrow A \frac{\cos(\theta)}{a^2} &= B a \cos(\theta) \end{aligned}$$

- (2) Discontinuity in the radial electric field at the surface must be related to the local surface charge density

$$\begin{aligned} \epsilon_0 (E_{out,r}|_{r=a} - E_{in,r}|_{r=a}) &= \sigma \\ \Rightarrow -\epsilon_0 \frac{\partial \phi_{out}(\vec{r})}{\partial r} \Big|_{r=a} + \epsilon_0 \frac{\partial \phi_{in}(\vec{r})}{\partial r} \Big|_{r=a} &= \sigma_0 \cos(\theta) \\ \Rightarrow 2A \frac{\cos(\theta)}{a^3} + B \cos(\theta) &= \frac{\sigma_0 \cos(\theta)}{\epsilon_0} \end{aligned}$$

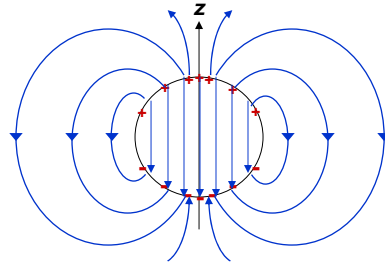
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A Non-Uniformly Charged Spherical Shell - IV

Solution:

$$\phi_{out}(\vec{r}) = \left(\frac{\sigma_0 a^3}{3\epsilon_0} \right) \frac{\cos(\theta)}{r^2}$$

$$\phi_{in}(\vec{r}) = \left(\frac{\sigma_0}{3\epsilon_0} \right) r \cos(\theta) = \left(\frac{\sigma_0}{3\epsilon_0} \right) z$$



A surface charge density of the form:

$$\sigma = \sigma_0 \cos(\theta)$$

produces a **uniform z-directed E-field** inside the charged shell and a **dipole-like E-field** outside the charged shell

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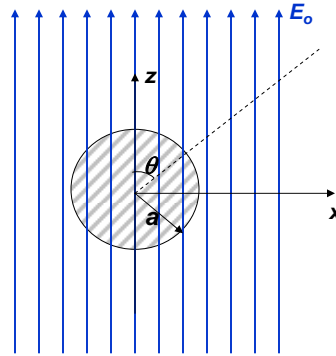
A Perfect Metallic Sphere in a Uniform E-Field - I

Consider a perfect metal sphere placed inside a constant and uniform z-directed E-field

The field lines shown on the right cannot be the right picture – there cannot be any E-field inside a perfect metal

Question: So what happens when a metal sphere is placed inside a constant and uniform electric field ??

Answer: Induced surface charges appear on the metal sphere that screen out the outside E-field



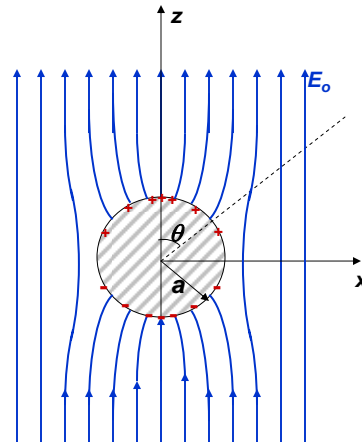
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A Perfect Metallic Sphere in a Uniform E-Field - II

Question: What kind of surface charge density gets induced on the metal sphere that produces a uniform E-field inside the sphere which completely cancels the applied E-field??

Answer: Induced surface charge density on the metal sphere must look like:

$$\sigma \propto \cos(\theta)$$



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A Perfect Metallic Sphere in a Uniform E-Field - III

Solution:

The outside potential must look like a combination of:

(i) The potential corresponding to a uniform z-directed E-field, and

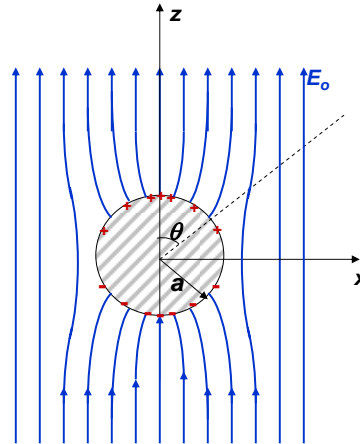
(ii) The dipole-like potential from the induced surface charge density on the metal sphere

So try the following solution for $r > a$:

$$\text{For } r > a: \phi_{out}(\vec{r}) = A \frac{\cos(\theta)}{r^2} - E_o r \cos(\theta)$$

For $r < a$: Metal sphere is equipotential with zero potential

$$\phi_{in}(\vec{r}) = 0$$



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A Perfect Metallic Sphere in a Uniform E-Field - IV

$$\text{For } r > a: \phi_{out}(\vec{r}) = A \frac{\cos(\theta)}{r^2} - E_o r \cos(\theta)$$

For $r < a$: Metal sphere is an equipotential with zero potential

$$\phi_{in}(\vec{r}) = 0$$

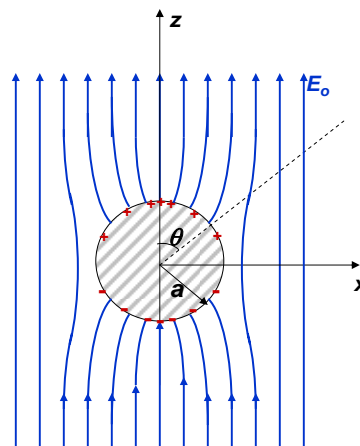
Boundary Condition:

(i) The potential at the surface of the metal must be continuous (i.e. the same inside and outside) and, therefore, must be zero

$$\begin{aligned} \phi_{in}(\vec{r})_{r=a} &= \phi_{out}(\vec{r})_{r=a} \\ \Rightarrow A \frac{\cos(\theta)}{a^2} - E_o a \cos(\theta) &= 0 \end{aligned}$$

Solution:

$$\phi_{out}(\vec{r}) = E_o a^3 \frac{\cos(\theta)}{r^2} - E_o r \cos(\theta)$$



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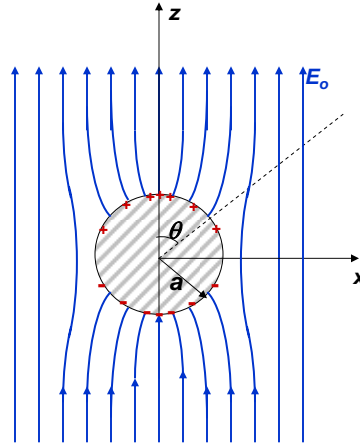
A Perfect Metallic Sphere in a Uniform E-Field - V

Surface Charge Density:

Now lets calculate the induced surface charge density on the metal sphere surface

$$\begin{aligned} \epsilon_0 (E_{out,r}|_{r=a} - E_{in,r}|_{r=a}) &= \sigma \\ \Rightarrow -\epsilon_0 \left. \frac{\partial \phi_{out}(\vec{r})}{\partial r} \right|_{r=a} - 0 &= \sigma \\ \Rightarrow \sigma &= 3 \epsilon_0 E_0 \cos(\theta) \end{aligned}$$

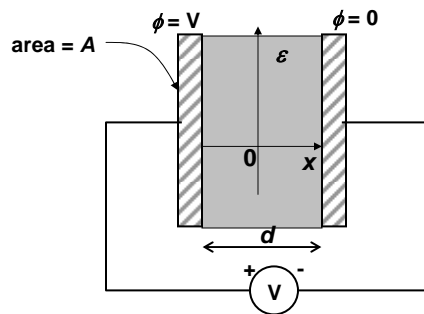
As expected, the induced surface charge density is just what is needed to completely cancel the applied E-field inside the metal sphere to give a net zero E-field



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Parallel Plate Capacitor with a Dielectric - I

Consider the parallel plate capacitor now with a dielectric between the plates



Need to find the potential $\phi(x)$ between the plates

Need to solve:

$$\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} = 0$$

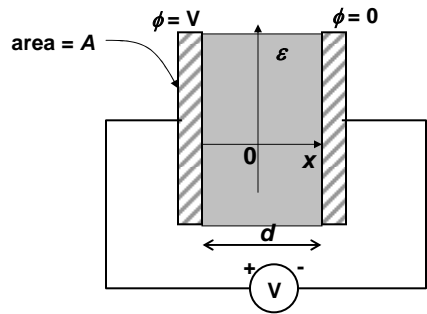
With the two boundary conditions:

$$\phi\left(x = -\frac{d}{2}\right) = V \quad \phi\left(x = +\frac{d}{2}\right) = 0$$

Solution is: $\phi(x) = V\left(\frac{1}{2} - \frac{x}{d}\right) \Rightarrow E_x(x) = -\frac{\partial \phi(x)}{\partial x} = \frac{V}{d}$

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Parallel Plate Capacitor with a Dielectric - II



Solution is:

$$\phi(x) = V \left(\frac{1}{2} - \frac{x}{d} \right)$$

$$\Rightarrow E_x(x) = -\frac{\partial \phi(x)}{\partial x} = \frac{V}{d}$$

Surface Charge Densities on Metal Plates

$$\epsilon_2 E_2 - \epsilon_1 E_1 = \sigma_u$$

On the left plate (σ_L)

- The field inside the metal is zero
- The field just outside the metal is V/d

$$\left(\epsilon \frac{V}{d} - 0 \right) = \sigma_L \Rightarrow \sigma_L = \epsilon \frac{V}{d}$$

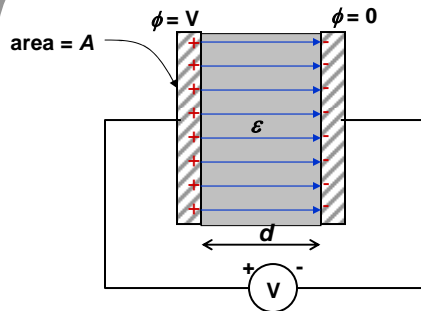
On the right plate (σ_R)

- The field just outside the metal is V/d
- The field inside the metal is zero

$$\left(0 - \epsilon \frac{V}{d} \right) = \sigma_R \Rightarrow \sigma_R = -\epsilon \frac{V}{d}$$

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Parallel Plate Capacitor with a Dielectric - III



Surface Charge Densities on Metal Plates

$$\sigma_L = \epsilon \frac{V}{d}$$

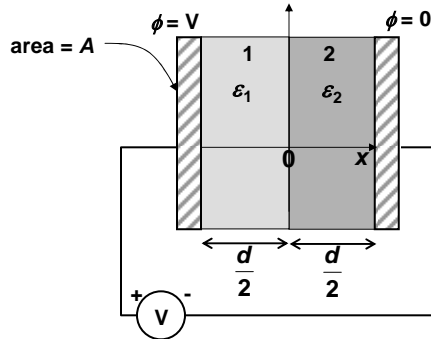
$$\sigma_R = -\epsilon \frac{V}{d}$$

Total charge on the positive plate: $Q = \frac{\epsilon A}{d} V$

Capacitance: $C = \frac{Q}{V} = \frac{dQ}{dV} = \frac{\epsilon A}{d}$

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Parallel Plate Capacitor with a Non-Uniform Dielectric - I



In region 1 need to solve:

$$\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} = 0$$

assume a solution:

$$\phi_1(x) = Ax + B$$

In region 2 need to solve:

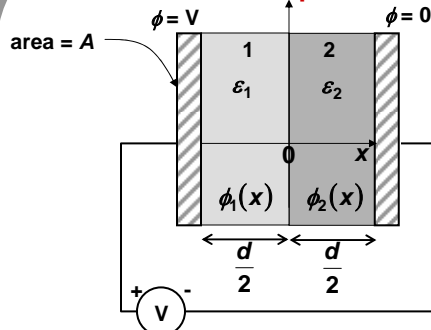
$$\nabla^2 \phi = \frac{\partial^2 \phi}{\partial x^2} = 0$$

assume a solution:

$$\phi_2(x) = Cx + D$$

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Parallel Plate Capacitor with a Non-Uniform Dielectric - II



In region 1:

$$\phi_1(x) = Ax + B$$

In region 2:

$$\phi_2(x) = Cx + D$$

Boundary Conditions: Need 4 boundary conditions to determine A, B, C, and D

$$(1) \phi_1\left(x = -\frac{d}{2}\right) = V \quad (2) \phi_2\left(x = +\frac{d}{2}\right) = 0$$

$$(3) \phi_1(x=0) = \phi_2(x=0) \quad \longrightarrow \quad \text{Continuity of the potential at } x=0$$

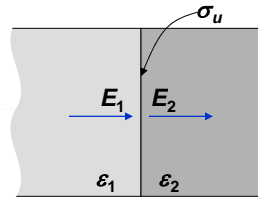
$$(4) \quad \varepsilon_2 E_2(x=0) - \varepsilon_1 E_1(x=0) = -\varepsilon_2 \left. \frac{\partial \phi_2}{\partial x} \right|_{x=0} + \varepsilon_1 \left. \frac{\partial \phi_1}{\partial x} \right|_{x=0} = \sigma_u = 0$$

Discontinuity of the normal component of the E-field at $x=0$

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Remember Your Boundary Conditions

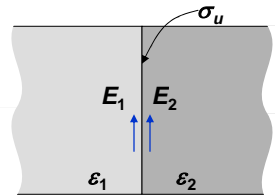
(1) The **discontinuity** of the **normal component** of the D-field at a material interface is related to the (unpaired) surface charge density at the interface



$$D_2 - D_1 = \sigma_u \quad \text{or} \quad \epsilon_2 E_2 - \epsilon_1 E_1 = \sigma_u$$

$$\text{or} \quad \epsilon_0 (E_2 - E_1) = \sigma_u + \sigma_p$$

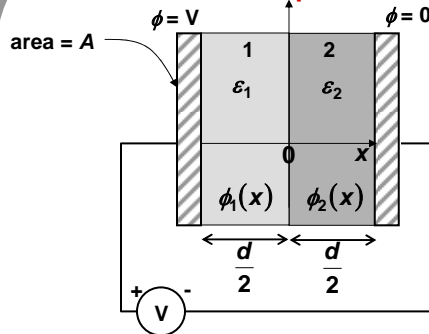
(2) The **parallel component** of the E-field at a material interface is always **continuous** at the interface (no change here)



$$(E_2 - E_1) = 0$$

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Parallel Plate Capacitor with a Non-Uniform Dielectric - III



In region 1:

$$\phi_1(x) = -V \frac{\epsilon_2}{\epsilon_1 + \epsilon_2} \left(\frac{2x}{d} - \frac{\epsilon_1}{\epsilon_2} \right)$$

In region 2:

$$\phi_2(x) = -V \frac{\epsilon_1}{\epsilon_1 + \epsilon_2} \left(\frac{2x}{d} - 1 \right)$$

Surface Charge Densities on Metal Plates
On the left plate (σ_L)

$$\sigma_L = \epsilon_1 E_x \left(x = -\frac{d}{2} \right) = \frac{V}{d} \frac{2 \epsilon_1 \epsilon_2}{\epsilon_1 + \epsilon_2}$$

On the right plate (σ_R)

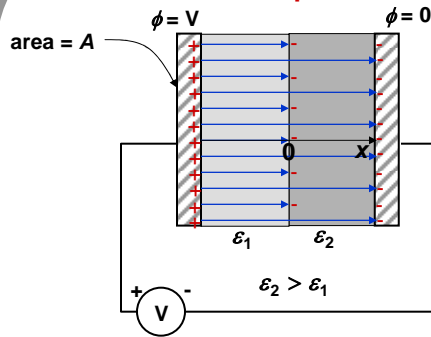
$$\sigma_R = -\epsilon_2 E_x \left(x = \frac{d}{2} \right) = -\frac{V}{d} \frac{2 \epsilon_1 \epsilon_2}{\epsilon_1 + \epsilon_2}$$

Capacitance

$$C = \frac{Q}{V} = \frac{\sigma_L A}{V} = \frac{2 \epsilon_1 \epsilon_2}{\epsilon_1 + \epsilon_2} \frac{A}{d}$$

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Parallel Plate Capacitor with a Non-Uniform Dielectric - IV



The discontinuity of the E-field at the interface of the two dielectrics is related to the surface charge density due to the paired charges (plot for $\epsilon_2 > \epsilon_1$)

Paired Surface Charge Density at the Interface of the Two Dielectrics:

$$\epsilon_0 (E_2 - E_1) = \sigma_u + \sigma_p$$

Use: $\epsilon_0 (E_2(x=0) - E_1(x=0)) = \sigma_u + \sigma_p = \sigma_p$

$$\sigma_p = -\epsilon_0 \frac{V}{d} \frac{2(\epsilon_2 - \epsilon_1)}{\epsilon_1 + \epsilon_2}$$

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