

Lecture 34

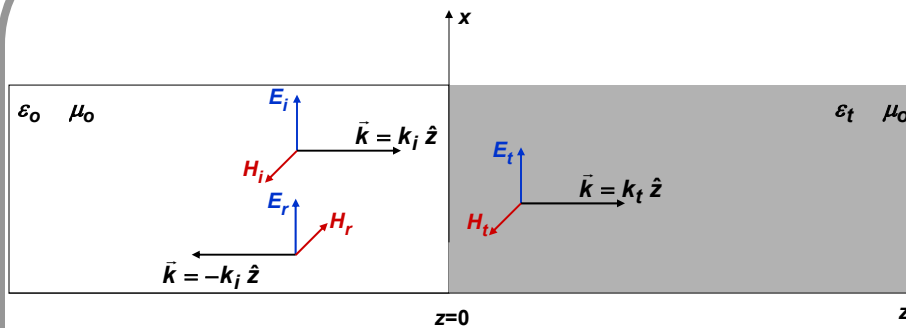
Electromagnetic Scattering

In this lecture you will learn:

- Scattering of electromagnetic waves from objects
- Rayleigh Scattering
- Why the sky is blue
- Radar range equation

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Scattering of Electromagnetic Waves from a Plane Interface

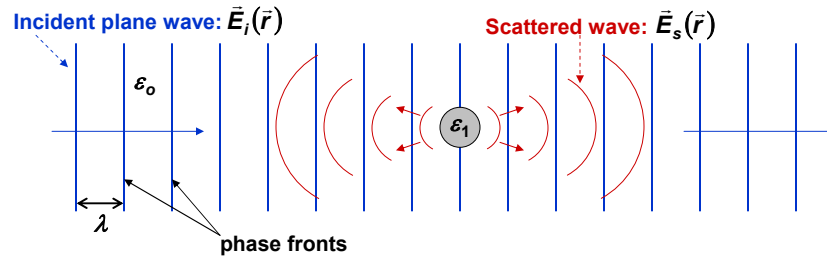


Incident, transmitted, and reflected waves are all plane waves

The reflected and transmitted waves can also be called the scattered waves

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Scattering of Electromagnetic Waves from Objects



Questions:

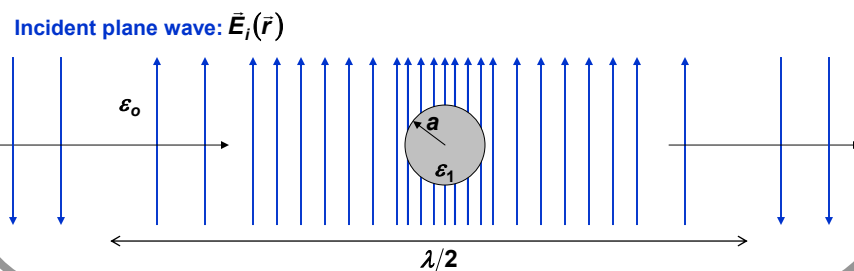
- How does one find the scattered field?
- How much power from the incident field goes into the scattered field?
- In which direction(s) does the scattered power go?

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Scattering of Electromagnetic Waves From Spherical Particles

Assumption:

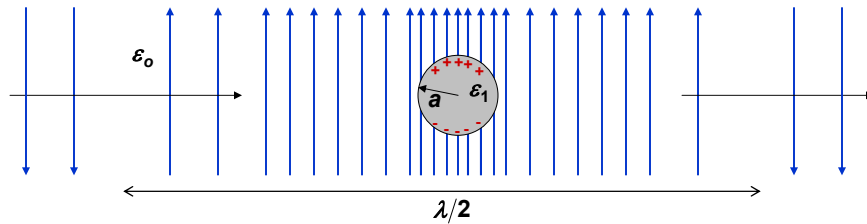
- Assume that the particle radius is much smaller than the wavelength of the incident wave, i.e.: $ka \ll 1$
 - When the above condition holds the scattering is called “Rayleigh Scattering”
 - When the above condition does not hold the scattering is called “Mie Scattering”
- When $ka \ll 1$, the particle sees a uniform E-field that is slowly oscillating in time



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Rayleigh Scattering: Basic Mechanism

Incident plane wave: $\vec{E}_i(\vec{r})$



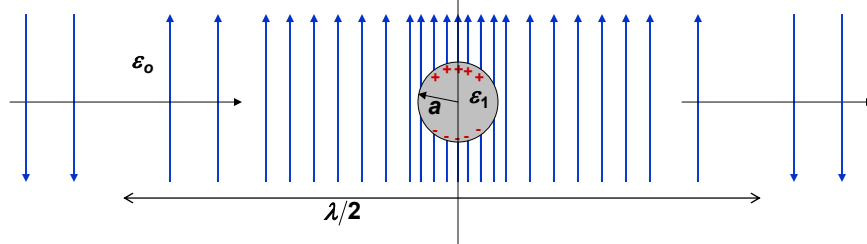
One way to understand scattering is as follows:

- i) The incident E-field induces a time-varying dipole moment in the sphere (recall the electrostatics problem of a dielectric sphere in a uniform E-field from homework 3)
- ii) The time-varying dipole radiates like a Hertzian dipole, and this radiation is the scattered radiation

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Rayleigh Scattering: Induced Dipole

Incident plane wave: $\vec{E}_i(\vec{r})$



Suppose the z-directed E-field phasor for the incident plane wave at the location of the particle is: $\vec{E}(\vec{r} = 0) = \hat{z} E_i$

From homework (3), the z-directed dipole moment ρ induced in a sphere in the presence of E-field E is:

$$\rho = 4\pi \epsilon_0 a^3 \left(\frac{\epsilon_1 - \epsilon_0}{\epsilon_1 + 2\epsilon_0} \right) E$$

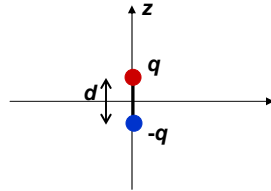
In the present case, the dipole moment phasor ρ induced in the sphere is therefore:

$$\rho = 4\pi \epsilon_0 a^3 \left(\frac{\epsilon_1 - \epsilon_0}{\epsilon_1 + 2\epsilon_0} \right) E_i$$

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Rayleigh Scattering: Scattered Radiation

Hertzian Dipole



$$\vec{E}_{ff}(\vec{r}) = \hat{\theta} \frac{j \eta_0 k l d}{4\pi r} \sin(\theta) e^{-jkr}$$

Dipole moment = $p = qd$

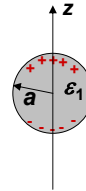
Current = $I = j \omega q$

$\Rightarrow ld = j \omega qd = j \omega p$

Therefore:

$$\vec{E}_{ff}(\vec{r}) = \hat{\theta} \frac{j \eta_0 k (j \omega p)}{4\pi r} \sin(\theta) e^{-jkr}$$

Induced Dipole

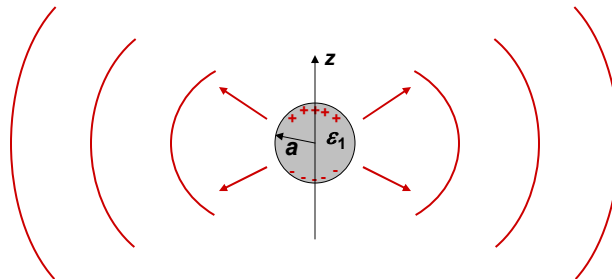


$$\text{Dipole moment} = p = 4\pi \epsilon_0 a^3 \left(\frac{\epsilon_1 - \epsilon_0}{\epsilon_1 + 2\epsilon_0} \right) E_i$$

The far-field scattered radiation is:

$$\begin{aligned} \vec{E}_{s-ff}(\vec{r}) &= \hat{\theta} \frac{j \eta_0 k (j \omega p)}{4\pi r} \sin(\theta) e^{-jkr} \\ &= -\hat{\theta} \frac{k^2 a^3}{r} \left(\frac{\epsilon_1 - \epsilon_0}{\epsilon_1 + 2\epsilon_0} \right) E_i \sin(\theta) e^{-jkr} \end{aligned}$$

Rayleigh Scattering: Total Scattered Power



$$\begin{aligned} P_s &= \int_0^{2\pi} \int_0^\pi \frac{|\vec{E}_{s-ff}(\vec{r})|^2}{2\eta_0} r^2 \sin(\theta) d\theta d\phi \\ &= \frac{4\pi}{3\eta_0} k^4 a^6 \left(\frac{\epsilon_1 - \epsilon_0}{\epsilon_1 + 2\epsilon_0} \right)^2 |E_i|^2 \end{aligned}$$

Rayleigh Scattering: Scattering Cross-Section

Total scattered power P_s from a dielectric sphere is:

$$P_s = \int_0^{2\pi} \int_0^\pi \frac{|\vec{E}_{s-ff}(\vec{r})|^2}{2\eta_0} r^2 \sin(\theta) d\theta d\phi = \frac{4\pi}{3\eta_0} k^4 a^6 \left(\frac{\epsilon_1 - \epsilon_0}{\epsilon_1 + 2\epsilon_0} \right)^2 |E_i|^2$$

The incident power per unit area was just the Poynting vector of the incident wave: $\frac{|\vec{E}_i(\vec{r})|^2}{2\eta_0}$

The **scattering cross-section** σ_s of a scatterer is defined as the area of a plane oriented perpendicular to the direction of incident wave that would intercept the same total incident power as the power P_s that the scatterer radiates

$$\sigma_s = \frac{P_s}{|\vec{E}_i(\vec{r})|^2 / 2\eta_0}$$

σ_s is also the ratio of the total scattered power to the power per unit area of the incident wave at the location of the scatterer

For the dielectric sphere:

$$\sigma_s = \frac{P_s}{|\vec{E}_i(\vec{r})|^2 / 2\eta_0} = \frac{8\pi}{3} k^4 a^6 \left(\frac{\epsilon_1 - \epsilon_0}{\epsilon_1 + 2\epsilon_0} \right)^2$$

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Rayleigh Scattering: Why is the Sky Blue

The scattering cross-section of a dielectric sphere is:

$$\sigma_s = \frac{8\pi}{3} k^4 a^6 \left(\frac{\epsilon_1 - \epsilon_0}{\epsilon_1 + 2\epsilon_0} \right)^2$$

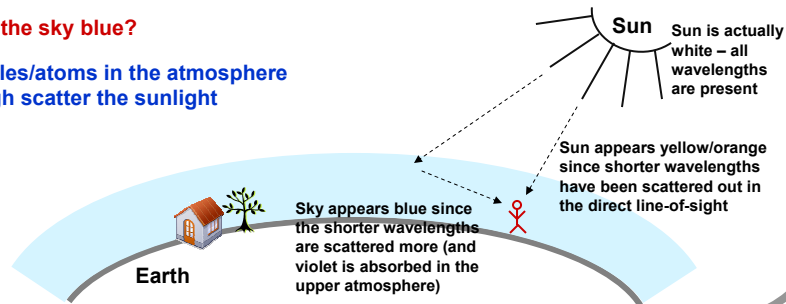
The scattered power is inversely proportional to the fourth power of the wavelength:

$$\sigma_s \propto k^4 \propto \frac{1}{\lambda^4}$$

Shorter wavelengths are scattered more than longer wavelengths in the Rayleigh limit

Why is the sky blue?

Molecules/atoms in the atmosphere Rayleigh scatter the sunlight



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Example: Radar Range Equation

Power per unit area at the location of the target:

$$S_{\text{target}} = \frac{P_{\text{in}}}{4\pi r^2} G(\theta, \phi)$$

If the scattering cross-section of the target is σ_s then the total scattered power P_s is:

$$P_s = \sigma_s S_{\text{target}} = \sigma_s \frac{P_{\text{in}}}{4\pi r^2} G(\theta, \phi)$$

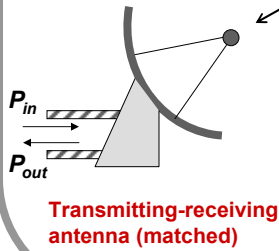
If the target scatters isotropically (equally in all directions) then the power P_{out} received by the matched antenna is:

$$P_{\text{out}} = \eta_p \frac{P_s}{4\pi r^2} A(\theta, \phi)$$

$$= \eta_p \frac{P_{\text{in}}}{(4\pi r^2)^2} \sigma_s G(\theta, \phi) A(\theta, \phi)$$

Radar range equation

$$\Rightarrow \frac{P_{\text{out}}}{P_{\text{in}}} = \eta_p \frac{\sigma_s G^2(\theta, \phi)}{(4\pi r^2)^2} \left(\frac{\lambda^2}{4\pi} \right) = \eta_p \frac{\sigma_s G^2(\theta, \phi) \lambda^2}{(4\pi)^3 r^4}$$



Transmitting-receiving antenna (matched)